Kendall coordinates are similar to Bookstein coordinates but location is removed in a different manner. We first need to define the Helmert sub-matrix which is used to remove location.

The Helmert sub-matrix $H$ is the $(k-1) \times k$ Helmert matrix without the first row. The full Helmert matrix $H^{F}$, which is commonly used in Statistics, is a square $k \times k$ orthogonal matrix with its first row of elements equal to $1 / \sqrt{ } k$, and the remaining rows are orthogonal to the first row. We drop the first row of $H^{F}$ so that the transformed $H X$ does not depend on the original location of the configuration.

Definition 2.5 The jth row of the Helmert sub-matrix $H$
is given by

$$
\begin{equation*}
\left(h_{j}, \ldots, h_{j},-j h_{j}, 0, \ldots, 0\right), \quad h_{j}=-\{j(j+1)\}^{-1 / 2} \tag{2.9}
\end{equation*}
$$

and so the jth row consists of $h_{j}$ repeated $j$ times, followed by $-j h_{j}$ and then $k-j-1$ zeros, $j=1, \ldots, k-1$.

For $k=3$ the full Helmert matrix is explicitly

$$
H^{F}=\left[\begin{array}{ccc}
1 / \sqrt{ } 3 & 1 / \sqrt{ } 3 & 1 / \sqrt{ } 3 \\
-1 / \sqrt{ } 2 & 1 / \sqrt{ } 2 & 0 \\
-1 / \sqrt{ } 6 & -1 / \sqrt{ } 6 & 2 / \sqrt{ } 6
\end{array}\right]
$$

and the Helmert sub-matrix is

$$
H=\left[\begin{array}{ccc}
-1 / \sqrt{ } 2 & 1 / \sqrt{ } 2 & 0 \\
-1 / \sqrt{ } 6 & -1 / \sqrt{ } 6 & 2 / \sqrt{ } 6
\end{array}\right] .
$$

For $k=4$ points the full Helmert matrix is

$$
H^{F}=\left[\begin{array}{cccc}
1 / 2 & 1 / 2 & 1 / 2 & 1 / 2 \\
-1 / \sqrt{ } 2 & 1 / \sqrt{ } 2 & 0 & 0 \\
-1 / \sqrt{ } 6 & -1 / \sqrt{ } 6 & 2 / \sqrt{ } 6 & 0 \\
-1 / \sqrt{ } 12 & -1 / \sqrt{ } 12 & -1 / \sqrt{ } 12 & 3 / \sqrt{ } 12
\end{array}\right]
$$

and the Helmert sub-matrix is

$$
H=\left[\begin{array}{cccc}
-1 / \sqrt{ } 2 & 1 / \sqrt{ } 2 & 0 & 0 \\
-1 / \sqrt{ } 6 & -1 / \sqrt{ } 6 & 2 / \sqrt{ } 6 & 0 \\
-1 / \sqrt{ } 12 & -1 / \sqrt{ } 12 & -1 / \sqrt{ } 12 & 3 / \sqrt{ } 12
\end{array}\right]
$$

Consider the original complex landmarks $z^{o}=\left(z_{1}^{o}, \ldots, z_{k}^{o}\right)^{\mathrm{T}}$ and remove location by pre-multiplying by the Helmert sub-matrix $H$ to give $z_{H}=H z^{o}=\left(z_{1}, \ldots, z_{k-1}\right)^{\mathrm{T}}$.

Definition 2.6 The Kendall coordinates are given by

$$
\begin{equation*}
u_{j}^{K}+i v_{j}^{K}=\frac{z_{j-1}}{z_{1}}, \quad j=3, \ldots, k . \tag{2.10}
\end{equation*}
$$

There is a simple 1-1 correspondence between Kendall
and Bookstein coordinates. If we write

$$
w^{B}=\left(u_{3}^{B}+i v_{3}^{B}, \ldots, u_{k}^{B}+i v_{k}^{B}\right)^{\mathrm{T}}
$$

for Bookstein coordinates and

$$
w^{K}=\left(u_{3}^{K}+i v_{3}^{K}, \ldots, u_{k}^{K}+i v_{k}^{K}\right)^{\mathrm{T}}
$$

for Kendall coordinates, then it follows that

$$
\begin{equation*}
w^{K}=\sqrt{ } 2 H_{1} w^{B} \tag{2.11}
\end{equation*}
$$

where $H_{1}$ is the lower right $(k-2) \times(k-2)$ partition matrix of the Helmert sub-matrix $H$. Note that $H_{1}^{\mathrm{T}} H_{1}=I_{k-2}-\frac{1}{k} 1_{k-2} 1_{k-2}^{\mathrm{T}}, \quad\left|H_{1}\right|^{2}=2 / k, \quad\left(H_{1}^{\mathrm{T}} H_{1}\right)^{-1}=I_{k-2}+\frac{1}{2} 1_{k-2} 1_{k-2}^{\mathrm{T}}$ so linear transformation from one coordinate system to the other is straightforward. The inverse transformation is

$$
w^{B}=\left(H_{1}^{\mathrm{T}} H_{1}\right)^{-1} H_{1}^{\mathrm{T}} w^{K} / \sqrt{2} .
$$

For $k=3$ we have the relationship

$$
u_{3}^{B}+i v_{3}^{B}=\frac{z_{3}^{o}-\frac{1}{2}\left(z_{1}^{o}+z_{2}^{o}\right)}{z_{2}^{o}-z_{1}^{o}}=\frac{\sqrt{ } 3}{2}\left(u_{3}^{K}+i v_{3}^{K}\right)
$$

and so Kendall coordinates in this case are the coordinates of the third landmark after transforming landmarks 1
and 2 to $(-1 / \sqrt{ } 3,0)$ and $(1 / \sqrt{ } 3,0)$ by the similarity transformations. The total transformation from $\left(z_{1}^{o}, z_{2}^{o}, z_{3}^{o}\right)$
to Kendall coordinates is

$$
z_{1}^{o} \rightarrow-\frac{1}{\sqrt{ } 3}, z_{2}^{o} \rightarrow \frac{1}{\sqrt{ } 3}, z_{3}^{o} \rightarrow u_{3}^{K}+i v_{3}^{K} .
$$

Throughout the text we shall often refer to the real $(2 k-4)$-vector of Kendall coordinates $u^{K}=$ $\left(u_{3}^{K}, \ldots, u_{k}^{K}, v_{3}^{K}, \ldots, v_{k}^{K}\right)^{\mathrm{T}}$, stacking the coordinates in this particular order.

### 2.3.5 Kendall's spherical coordinates for triangles

For $k=3$ we will see in Section 4.2.4 that the shape space is a sphere with radius $\frac{1}{2}$. A mapping from Kendall coordinates to the sphere of radius $\frac{1}{2}$ is

$$
\begin{equation*}
x=\frac{1-r^{2}}{2\left(1+r^{2}\right)}, \quad y=\frac{u_{3}^{K}}{1+r^{2}}, \quad z=\frac{v_{3}^{K}}{1+r^{2}} \tag{2.12}
\end{equation*}
$$

and $r^{2}=\left(u_{3}^{K}\right)^{2}+\left(v_{3}^{K}\right)^{2}$, so that $x^{2}+y^{2}+z^{2}=\frac{1}{4}$, where $u_{3}^{K}$ and $v_{3}^{K}$ are Kendall coordinates of Section 2.3.4.

Definition 2.7 Kendall's spherical coordinates $(\theta, \phi)$ are
given by the polar coordinates
$\frac{1}{2} \sin \theta \cos \phi=\frac{1-r^{2}}{2\left(1+r^{2}\right)}, \frac{1}{2} \sin \theta \sin \phi=\frac{u_{3}^{K}}{1+r^{2}}, \quad \frac{1}{2} \cos \theta=\frac{v_{3}^{K}}{1+r^{2}}$,
where $0 \leq \theta \leq \pi$ is the angle of latitude and $0 \leq \phi<2 \pi$ is the angle of longitude.

The relationship between $\left(u_{3}^{K}, v_{3}^{K}\right)$ and the spherical shape variables (Mardia, 1989c) is given by

$$
\begin{align*}
u_{3}^{K} & =\frac{\sin \theta \sin \phi}{1+\sin \theta \cos \phi} \\
v_{3}^{K} & =\frac{\cos \theta}{1+\sin \theta \cos \phi} \tag{2.14}
\end{align*}
$$

The sphere can be partitioned into 6 lunes and 12 half-
lunes. In order to make the terminology clear, one example
full-lune is $0 \leq \phi \leq \pi / 3,0 \leq \theta \leq \pi$ and one example
half-lune is $0 \leq \phi \leq \pi / 3,0 \leq \theta \leq \pi / 2$.
In Figure 27 we see triangle shapes located on the spherical shape space. The equilateral triangle with anticlockwise labelling corresponds to the 'North pole' $(\theta=$ 0 ) and the reflected equilateral triangle (with clockwise labelling) is at the 'South pole' $(\theta=\pi)$. The flat triangles (three collinear points) lie around the equator $(\theta=\pi / 2)$. The isosceles triangles lie on the meridians $\phi=$ $0, \pi / 3,2 \pi / 3, \pi, 4 \pi / 3,5 \pi / 3$. The right-angled triangles lie on three small circles given by

$$
\sin \theta \cos \left(\phi-\frac{2 k \pi}{3}\right)=\frac{1}{2}, k=0,1,2
$$

and we see the arc of unlabelled right-angled triangles on the front half-lune in Figure 27.

Reflections of triangles in the upper hemisphere at $(\theta, \phi)$ are located in the lower hemisphere at $(\pi-\theta, \phi)$. In addition, permuting the triangle labels gives rise to points in each of the six equal half-lunes in each hemisphere.

Thus, if invariance under labelling and reflection was required, then we would be restricted to one of these halflunes, for example the sphere surface defined by $0 \leq \phi \leq$ $\pi / 3,0 \leq \theta \leq \pi / 2$. Consider a triangle with labels A,B and C , and edge lengths $\mathrm{AB}, \mathrm{BC}$ and AC . If the labelling and reflection of the points was unimportant, then we could relabel each triangle so that, for example, $\mathrm{AB} \geq \mathrm{AC} \geq \mathrm{BC}$ and point C is above the baseline AB .

For practical analysis and the presentation of data it is often desirable to use a suitable projection of the sphere for triangle shapes. Kendall (1983) defined an equal area


Reflected equilateral (South pole)

Figure 27 Kendall's spherical shape space for triangles in $m=2$ dimensions. The shape coordinates are the latitude $\theta$ (with zero at the North pole) and the longitude $\phi$.
projection of one of the half-lunes of the shape sphere to display unlabelled triangle shapes. The projected lune is bell-shaped and this graphical tool is also known as 'Kendall's Bell' or the spherical blackboard (an example is given later in Figure 136).

An alternative equal-area projection is the Schmidt net (Mardia, 1989c) otherwise known as the Lambert projection given by

$$
\xi=2 \sin \left(\frac{\theta}{2}\right), \psi=\phi ; 0 \leq \xi \leq \sqrt{ } 2,0 \leq \psi<2 \pi
$$

In Figure 28 we see a plot of one of the half-lunes on the upper hemisphere of shape space projected onto the Schmidt net. Example triangles are drawn with their centroids at polar coordinates $(\xi, \psi)$ in the Schmidt net.


Figure 28 Part of the shape space of triangles projected onto the equal-area projection Schmidt net. If relabelling and reflection was not important, then all triangles could be projected into this sector.

